

# Tutorial on Intermittency and the multifractal formalism

Edinburgh, LMS-EPSRC Short Course

Following previous works [6, 7], the (Eulerian) velocity increment  $\delta_\ell u(x) = u(x + \ell) - u(x)$  is modeled as the product of two **independent** random variables (the scale  $\ell$  is assumed smaller than the integral length scale  $L$ ):

$$\delta_\ell u \stackrel{\text{law}}{=} \beta_\ell \times \delta ,$$

where  $\delta$  is a Gaussian zero-average unit-variance random variable, and  $\beta_\ell$  a *stochastic variance*, namely a positive random variable, assumed of the form

$$\beta_\ell = \sigma \left( \frac{\ell}{L} \right)^h ,$$

with  $\sigma^2 = \langle (\delta_L u)^2 \rangle$  the large-scale velocity increment variance and  $h$  a random exponent (called the Hölder exponent in multifractal terminology [6]) of density

$$\mathcal{P}_h^{(\ell)}(h) = \frac{1}{\mathcal{Z}(\ell)} \left( \frac{\ell}{L} \right)^{1-\mathcal{D}(h)} ,$$

where  $\mathcal{D}(h)$  is a (scale independent) parameter function called the singularity spectrum and  $\mathcal{Z}(\ell)$  a normalizing function.

- Compute the  $q^{\text{th}}$ -order structure function  $M_q(\ell) = \langle |\delta_\ell u|^q \rangle$ .  
hint:  $\langle |\delta|^p \rangle = \Gamma\left(\frac{p+1}{2}\right) / \sqrt{2p\pi}$  where  $\Gamma$  is the Gamma function.
- assuming that  $\min_h [1 - \mathcal{D}(h)] = 0$  and considering vanishing scales, i.e.  $\ell/L \rightarrow 0$ , show that structure functions behave as power-laws,

$$M_q(\ell) \sim \left( \frac{\ell}{L} \right)^{\zeta_q} ,$$

with

$$\zeta_q = \min_h [qh + 1 - \mathcal{D}(h)] .$$

hint: apply a steepest-descent approximation.

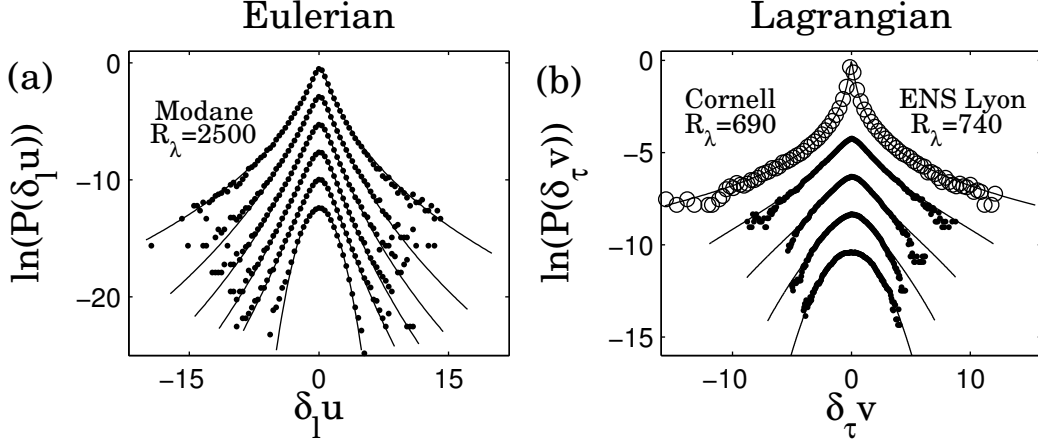


Figure 1: (a) PDFs of signed longitudinal velocity increments in Modane data [1]. Represented scales (from top to bottom):  $\ln(\ell/L) = -6.4137, -5.6028, -4.6645, -3.6411, -2.7501, -1.8598, -0.8685, 0.1226$ . All curves are arbitrarily vertically shifted for the sake of clarity and their variance is set to unity. The solid curves correspond to theoretical predictions (see Ref. [2]). (b) PDFs of Lagrangian temporal increments from the ENS-Lyon experiment, for time lags (from bottom to top, symbols  $\bullet$ )  $\tau/T = 0.07, 0.16, 0.35, 1$  and from Cornell acceleration data (symbols  $\circ$ , from Ref. [4]). Also, the curves are displayed with an arbitrary vertical shift for clarity, the variance is set to unity at any scales, and the original axis for the acceleration PDF ( $\circ$ ) has been shrunk by a factor 4. Solid lines correspond to theoretical predictions (see Ref. [5]).

- Lognormal model:

$$\mathcal{D}^{\text{LN}}(h) = 1 - \frac{(h - c_1)^2}{2c_2},$$

the parameter  $c_2$  is called the intermittency coefficient. Compute  $\zeta_q^{\text{LN}}$ .

- She-Lévéque model:

$$\mathcal{D}^{\text{SL}}(h) = -1 + 3 \left[ \frac{1 + \ln(\ln(3/2))}{\ln(3/2)} - 1 \right] (h-1/9) - \frac{3}{\ln(3/2)} (h-1/9) \ln(h-1/9).$$

Compute  $\zeta_q^{\text{SL}}$ .

- Finally, show that the velocity increment probability density function in the general case reads:

$$\mathcal{P}_{\delta_{\ell}u}(\delta_{\ell}u) = \int_{h_{\min}}^{h_{\max}} \frac{1}{\sigma} \left(\frac{\ell}{L}\right)^{-h} \mathcal{P}_{\delta} \left[ \frac{\delta_{\ell}u}{\sigma} \left(\frac{\ell}{L}\right)^{-h} \right] \mathcal{P}_h^{(\ell)}(h) dh ,$$

where  $\mathcal{P}_{\delta}(x) = \exp(-x^2/2)/\sqrt{2\pi}$  is the PDF of a unit-variance zero-mean Gaussian variable.

## References

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