

Low Energy Constants in QCD from Matrix Models

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Engineering and Physical Sciences
Research Council

& ENRAGE

collaboration with:

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Outline & Motivation

1. **QCD** → epsilon regime chiral Perturbation theory:

$$\text{box } V = L^4$$

$$\Lambda_{\text{QCD}} \gg p \sim 1/L \gg m_\pi \sim 1/L^2$$

LEC	F Pion decay constant	Σ chiral condensate
source term	μ chemical potential	m quark mass

2. Equivalent Matrix Models: easier

3. Lattice Data: fitting plots

• Results for:

spectral density & individual \not{D} -eigenvalues
currents $\langle S(x)S(0) \rangle$

Q**C**D and Banks-Casher Relation

$$\mathcal{Z}_{\text{QCD}} = \int [DA][d\Psi] \exp \left[- \int \bar{\Psi} (\not{D}[A] + M) \Psi - S_{YM}[A] + i\Theta \int F \tilde{F} \right]$$

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$$\text{in box } V : = \sum_{\nu} \int \prod_i^V [d\lambda_i[A]] e^{+i\Theta\nu} \prod_f^{N_f} m_f^\nu [\lambda_i^2 + m_f^2] e^{-S_{YM}[\{\lambda_j\}]}$$

$$\text{spectral density } \rho_{\not{D}}(\lambda) \equiv \langle \sum'_i \delta(\lambda - \lambda_i) \rangle_{\text{QCD}}$$

$$\boxed{\text{Banks-Casher} \quad \rho_{\not{D}}(\lambda \approx 0) = \frac{1}{\pi} V \Sigma, \quad \Sigma = \langle \bar{\Psi} \Psi \rangle}$$

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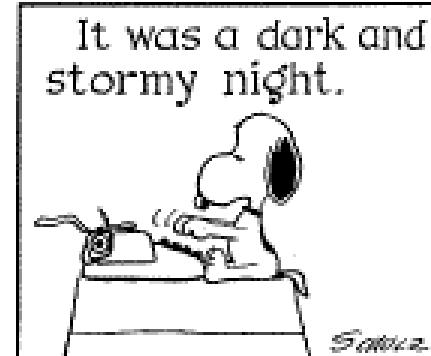
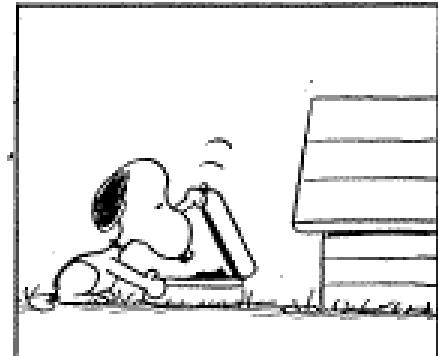
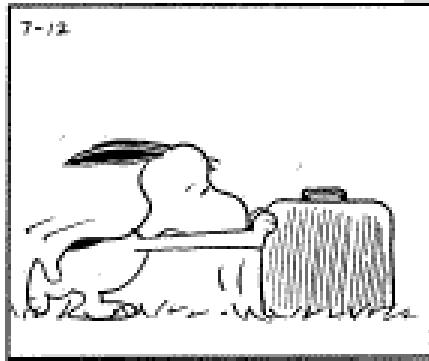
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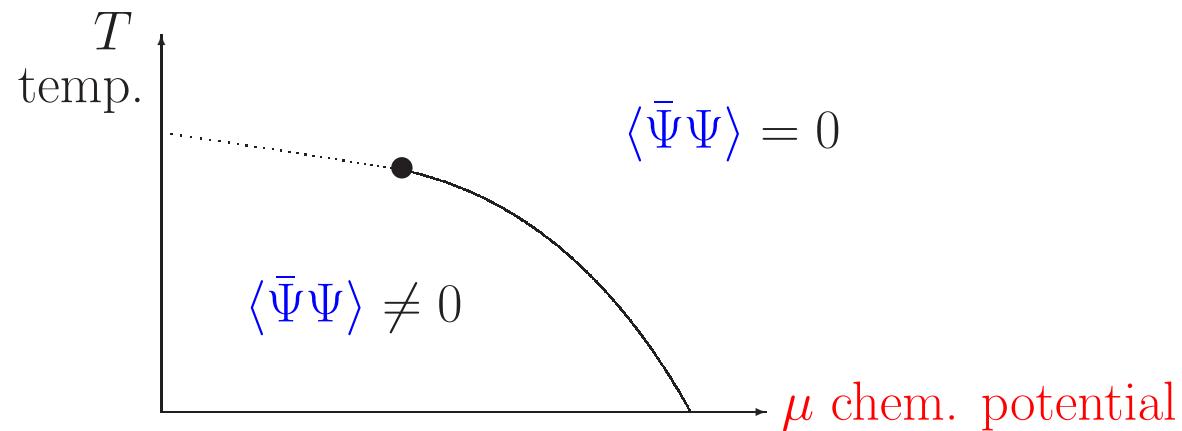
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Cartoon of QCD Phasediagram

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schematic phase diagram for **QCD** with 2 light flavors [Halasz et al. 98]

Matrix Models

$$\mathcal{Z}_{MM} \equiv \int dW \prod_{f=1}^{N_f} \det \begin{pmatrix} m_f & iW \\ iW^\dagger & m_f \end{pmatrix} e^{-N\sigma^2 \text{Tr } WW^\dagger}$$

[Shuryak, Verbaarschot 93]

fixed topology: $W = N \times (N + \nu)$ random matrix

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[Stephanov 96]

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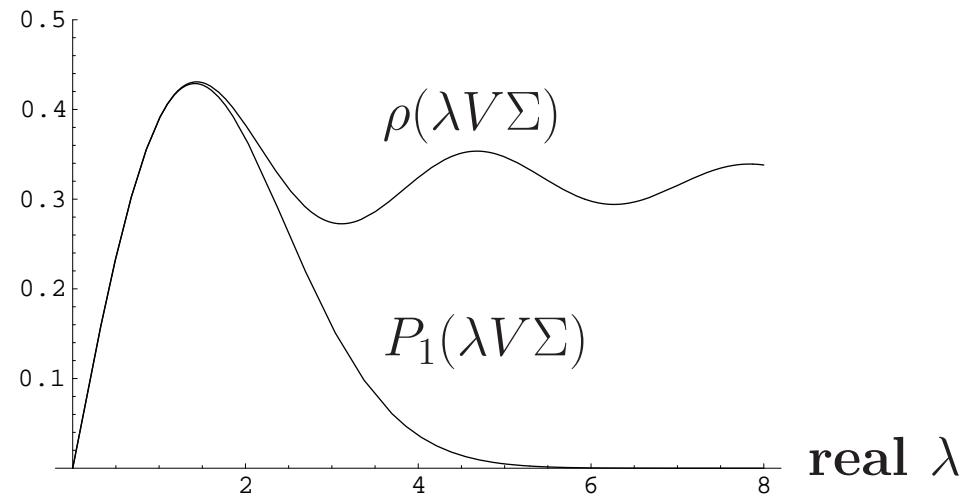
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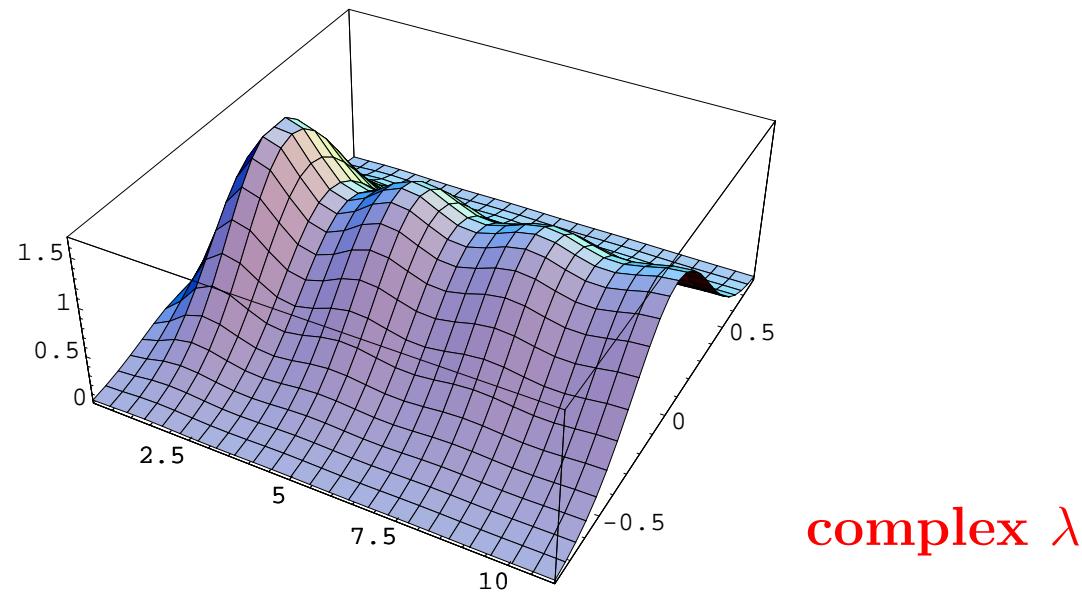
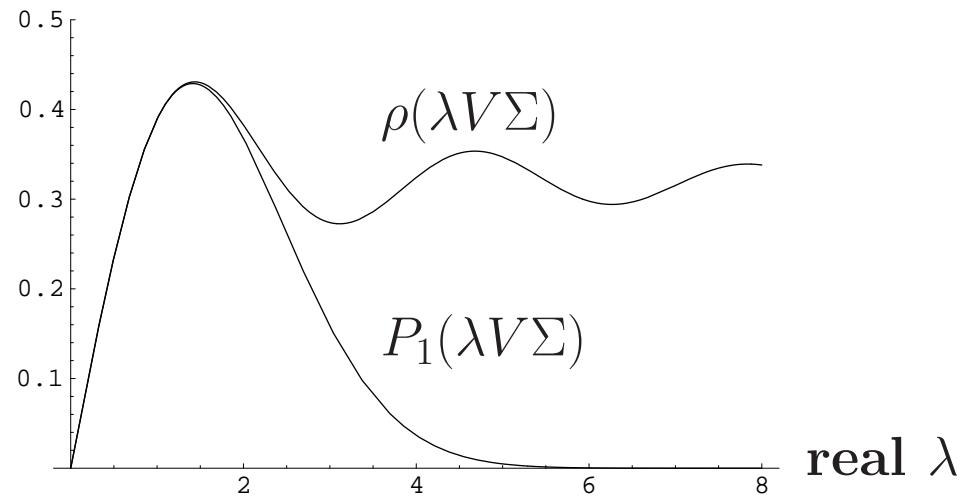
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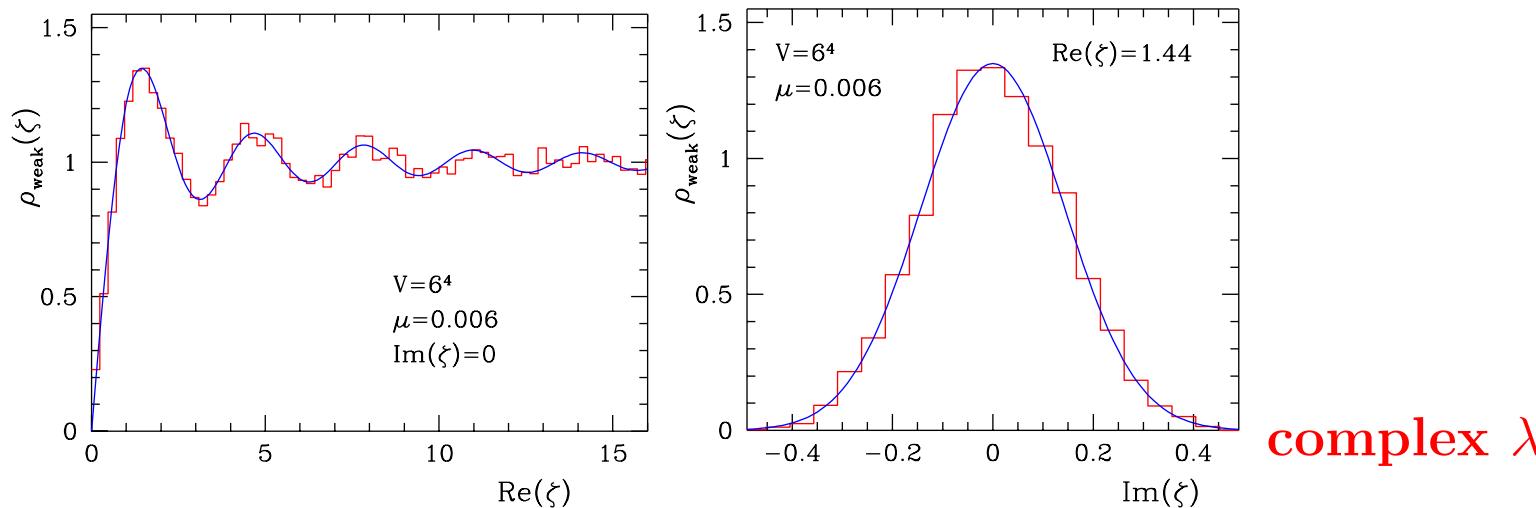
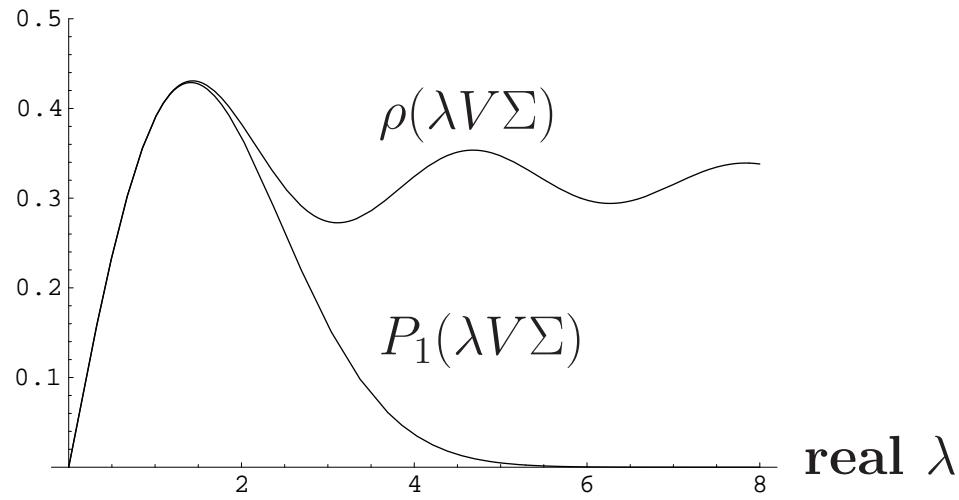
Spectral Density and Individual Eigenvalues



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Chiral Perturbation Theory and epsilon-Regime

spontaneous χ SB: $SU(N_f) \times SU(N_f) \longrightarrow SU(N_f)$

$$\mathcal{L} = \frac{1}{4} F^2 \nabla_a \mathbf{U}(x) \nabla_a \mathbf{U}^\dagger(x) + \frac{1}{2} \Sigma M \left(e^{i \frac{\Theta}{N_f}} \mathbf{U}(x) + e^{-i \frac{\Theta}{N_f}} \mathbf{U}^\dagger(x) \right)$$

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- zero-momentum modes non-perturbativ:

$$\mathbf{U}(x) = \mathbf{U}_0 e^{i\xi(x)}$$

$$LO : \mathcal{Z}_{\varepsilon\chi PT} = \int d\mathbf{U}_0 \det[\mathbf{U}_0]^\nu e^{-\frac{1}{2}\Sigma V \text{Tr}(\mathbf{M}(\mathbf{U}_0 + \mathbf{U}_0^\dagger))} \times \int [d\xi(x)] e^{-\int \frac{1}{2} \partial_a \xi \partial_a \xi}$$

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Objects of study: Currents & Densities

- **currents:** scalar source $S(x) = \bar{\Psi}(x)\mathbf{1}_{N_f}\Psi(x) \leftrightarrow \text{mass}$
[Hansen 90]

$$\langle S(x)S(0) \rangle_{\varepsilon\chi PT} \sim \frac{1}{Z_{\varepsilon\chi PT}} \partial_m^2 Z_{\varepsilon\chi PT} = \left\langle \left(\text{Tr}[U_0 + U_0^\dagger] \right)^2 \right\rangle_{\varepsilon\chi PT}$$

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- **densities:** SUSY- $\varepsilon\chi PT$ [Damgaard et al 98, Basile, G.A. 07]

$$\Sigma_\nu(m) \equiv \partial_m \left\langle \frac{\det[\not{D} + m]}{\det[\not{D} + m']} \right\rangle_{\varepsilon\chi PT}^{m=m'} = \left\langle \text{Tr} \frac{1}{\not{D} + m} \right\rangle_{\varepsilon\chi PT} = \int d\lambda \frac{\rho_{\not{D}}(\lambda) 2m}{\lambda^2 + m^2}$$

$\rho_{\not{D}}(m) = \Im m[\Sigma_\nu(m)] \quad (\text{or } \partial_m^* \Sigma_\nu(m) \text{ when } \mu \neq 0)$

1st eigenvalue : $P_1(\lambda) = \partial_\lambda \sum_k \frac{(-)^k}{k!} \int_0^\lambda d\lambda_1 \dots d\lambda_k \rho_k(\lambda_1, \dots, \lambda_k)$

[G.A., Damgaard 04]

Results: Some Examples

scaling $\hat{\lambda} = V \Sigma \lambda$

- **density:** $\rho(\hat{\lambda}) = \int_0^1 dt J_\nu(t\hat{\lambda}) J_\nu(t\hat{\lambda})$
- **k -point:** $\rho_k(\{\hat{\lambda}\}) = \det_{i,j} [\hat{\lambda}_i \hat{\lambda}_j] / \text{Vandermonde}$

$$\Rightarrow P_1(\hat{\lambda}) = \partial_{\hat{\lambda}} \sum_k \int \rho_k = \frac{1}{2} \hat{\lambda} \exp[\frac{1}{4} \hat{\lambda}^2]$$

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scaling $\hat{\lambda} = V\Sigma\lambda, \hat{\mu}^2 = VF^2\mu^2$

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- **current:**

$$\langle S(x)S(0) \rangle_{\varepsilon\chi PT} = \partial_{\hat{m}}^2 \log[\mathcal{Z}_{eff}] + \bar{\Delta}(x)(A + \hat{\mu}^2 B + \hat{\mu}^4 C)$$

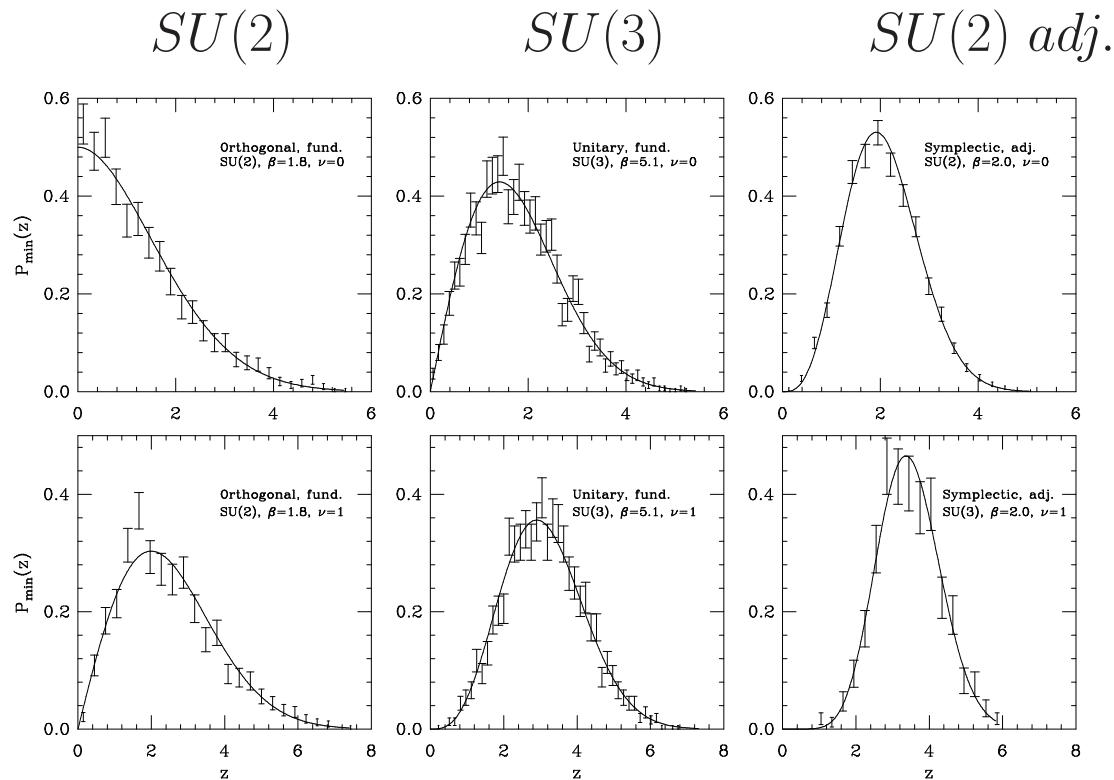
$$\boxed{\Sigma_{eff} = \Sigma(1 + \frac{c_1}{\sqrt{V}}) \quad F_{eff} = F(1 + \frac{c_2}{\sqrt{V}}), \quad \text{and } \Delta \sim \frac{1}{\sqrt{V}}}$$

$$\text{NEW } \mathcal{Z}_{\varepsilon\chi PT} = \int dU_0 \det[U_0]^\nu e^{-\text{Tr}\left(\frac{1}{2}\Sigma VM(\textcolor{violet}{U}_0 + U_0^\dagger) + \sum_K a_K ([B, U_0][B, U_0^\dagger])^K\right)}$$

Real Spectra vs. Lattice

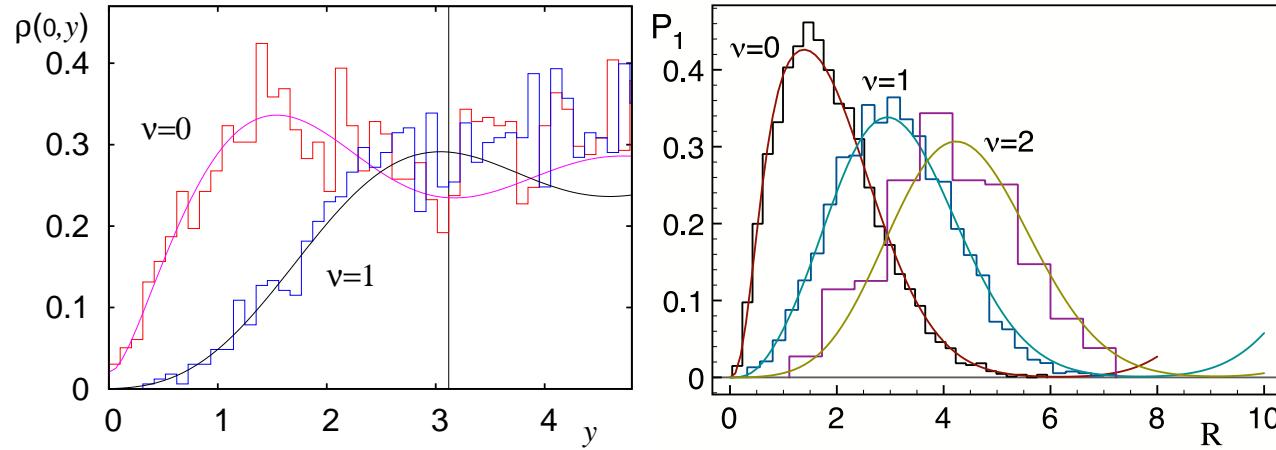
[R.G. Edwards et al 99]

$\nu = 0$



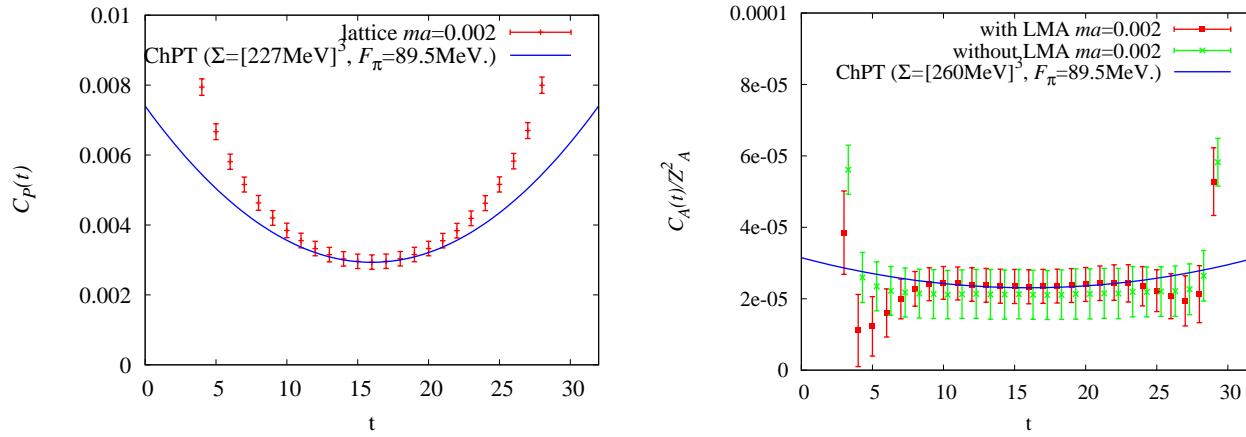
$\nu = 1$

Complex Spectra vs. Lattice



density ρ (left) [Bloch, Wettig 06], P_k (right) [+ G.A., Shifrin 07]

Currents vs. Lattice



$N_f = 2$ Pseudo-scalar (left) and Axial-vector (right)

[JLQCD 07]

Some Open Problems

- complete all χ SB classes with $\mu \neq 0$:
real and complex spectrum of adjoint **QCD**
- simple form for $P_1(\lambda)$ on \mathbb{C} ?
- Polyakov loop \longleftrightarrow $\not D$ eigenvalues
- other applications of new group integral ?