

# Fractons

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# Plan of talk

3D Plaquette Ising (Gonihedric) Model

Toric Code (and anyons)

Toric Code from gauging a global symmetry

X-Cube model (and fractons)

X-Cube model from gauging a subsystem symmetry

# Plaquette Ising/Gonihedric model

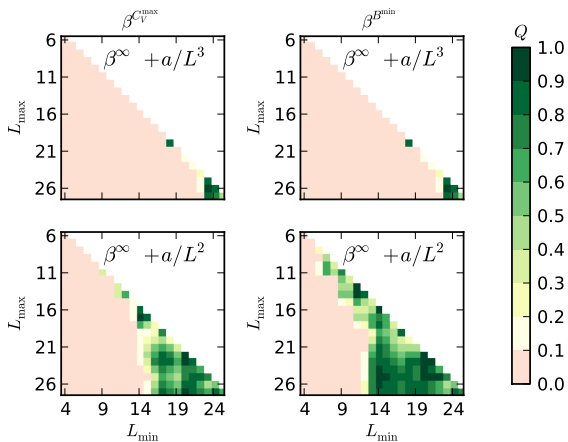
Hamiltonian

$$\mathcal{H} = - \sum_{\square} \sigma_i \sigma_j \sigma_k \sigma_l$$

Spins at vertices of 3D cubic lattice

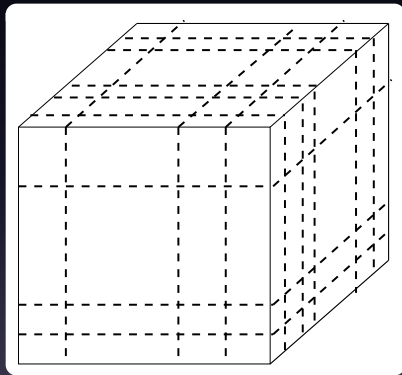
Strong first order phase transition

# Scaling oddity (Janke/Mueller)



Standard  $1/L^3$  gives much poorer quality

# Typical Ground state



# 1st Order FSS with Exponential Degeneracy

Normally  $q$  is constant

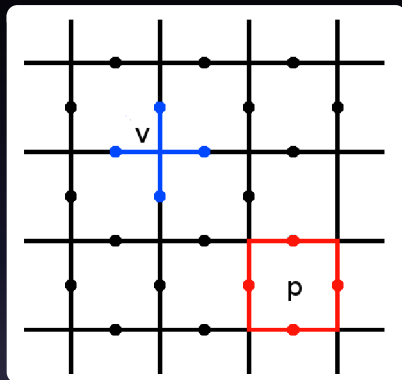
If instead  $q \propto e^L$  ( $q = e^{(3 \ln 2)L}$ ), as in Gonihedric model

$$\beta^{C_v^{\max}}(L) = \beta^\infty - \frac{\ln q}{L^d \Delta e} + \dots$$

becomes

$$\beta^{C_v^{\max}}(L) = \beta^\infty - \frac{3 \ln 2}{L^{d-1} \Delta e} + \dots$$

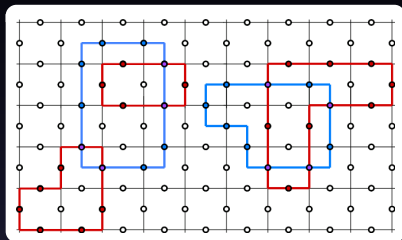
# Toric Code



$$A_v = \prod_{i \in v} \tau_i^x, \quad B_p = \prod_{i \in p} \tau_i^z$$

$$H = -J_v \sum_v A_v - J_p \sum_p B_p$$

# Toric Code: Ground State

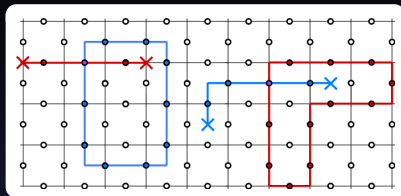


$$|\xi_0\rangle = \prod_v \frac{1}{\sqrt{2}} (\mathbb{1}_v + A_v) \underbrace{|0\rangle \otimes \dots \otimes |0\rangle}_{N_e \text{ times}}$$

A “Loop Soup”



# Toric Code: Excitations



Defects (i.e. quasiparticles) appear on the end of strings

$$W_e = \prod \tau_z. \quad W_m = \prod \tau_x$$

Braiding excitations reveals anyonic behavior

# Toric Code: Anyons

$e, m$  bosonic w.r.t. themselves

Take  $e$  for a walk around  $m$ , gives  $-1$  phase  $\implies$  anyons

Other interesting properties, topological degeneracy of ground state etc

# Gauging: From here to there:

**Here** - Quantum Transverse Ising:

$$\mathcal{H} = -\beta \sum_{\langle i,j \rangle} \sigma_i^z \sigma_j^z - h \sum_i \sigma_i^x$$

Gauge the global  $\mathbb{Z}_2$  symmetry

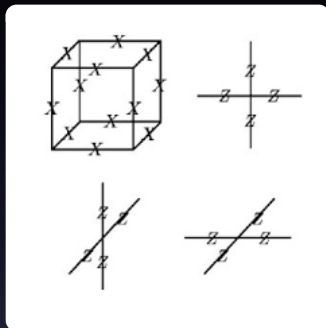
$$\mathcal{H} = -\beta \sum_{\langle i,j \rangle} \sigma_i^z \tau_{ij}^z \sigma_j^z - h \sum_i \sigma_i^x - J_p \sum_{\square} \tau_i^z \tau_j^z \tau_k^z \tau_l^z$$

$\beta \rightarrow 0$ , gauge invariance:  $\sigma_i^x \prod_{i \in \nu} \tau_i^x = 1$

**There** - Toric Code:

$$\mathcal{H} = -h \sum_{\nu} A_{\nu} - J_p \sum_p B_p$$

# The X-cube Model

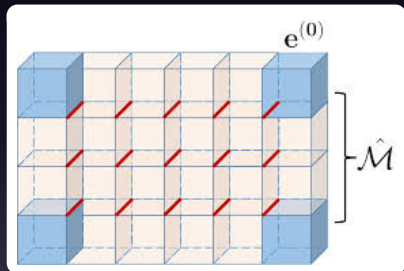


$$A_c = \prod_{i \in \square} \tau_i^x, \quad B_i^{xy,yz,xz} = \prod_{j \in +,i} \tau_j^z$$

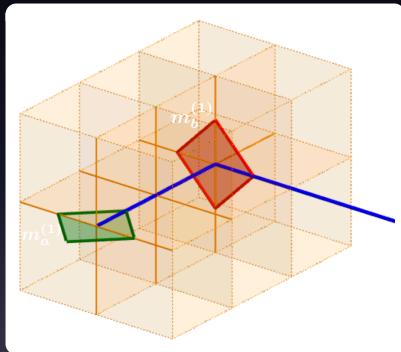
$$H = -J_{\square} \sum A_{\square} - J_{xy} \sum_i B_i^{xy} - J_{yz} \sum_i B_i^{yz} - J_{xz} \sum_i B_i^{xz}$$

# Toric Code: Fractons

Electric excitations  $\tau_Z$



Magnetic excitations  $\tau_X$



Pics c/o Vijay et.al.

# From here to there: Subsystem Symmetry Gauging

Here - Quantum Transverse Plaquette Ising:

$$\mathcal{H} = -\beta \sum_{\square} \sigma_i^z \sigma_j^z \sigma_k^z \sigma_l^z - h \sum_i \sigma_i^x$$

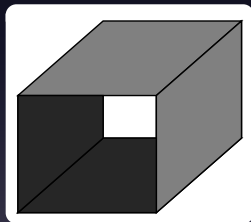
Gauge the  $\mathbb{Z}_2$  subsystem symmetry

$$\mathcal{H} = -\beta \sum_{\square} \tau_{\square}^z \sigma_i^z \sigma_j^z \sigma_k^z \sigma_l^z - h \sum_i \sigma_i^x + \dots$$

A picture, or two, is worth a thousand words

# From here to there: Subsystem Symmetry Gauging II

Equivalent of plaquette flux term in 2D is matchbox (not cube)



Gives  $B_i^{xy,yz,xz} = \prod_{j \in +,i} \sigma_j^z$  flux terms

# From here to there: Subsystem Symmetry Gauging

There (almost)

$$\mathcal{H} = -\beta \sum_{\square} \tau_{\square}^z \sigma_i^z \sigma_j^z \sigma_k^z \sigma_l^z - h \sum_i \sigma_i^x - J_{xy} \sum_i B_i^{xy} - J_{yz} \sum_i B_i^{yz} - J_{xz} \sum_i B_i^{xz}$$

$\beta \rightarrow 0$ , gauge invariance:  $\sigma_i^x \prod_i \tau_i^x = 1$

There

$$H = -h \sum_{\square} A_{\square} - J_{xy} \sum_i B_i^{xy} - J_{yz} \sum_i B_i^{yz} - J_{xz} \sum_i B_i^{xz}$$



## 2 second, 2 line summary

Gauge global  $\mathbb{Z}_2$ , get Toric Code, anyons etc

Gauge subsystem symmetry, get fractons  
(reduced mobility anyons)

# References

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